

Skyrmions in Condensed Matter Systems

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Skyrmion Anniversary, DAMTP

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Outline

Introduction

Solid State Systems

Quantum Hall Ferromagnets

Chiral Magnets

Atomic BECs

Summary

Introduction: Skyrmions

Skyrme model [T. H. R. Skyrme, Proc. Roy. Soc. **260**, 127 (1961)]

$\vec{n}(\mathbf{r}) = (n_1, n_2, n_3, n_4)$ with $|\vec{n}|^2 = 1$
in 3D, with $\vec{n}(\mathbf{r} \rightarrow \infty) = \vec{n}_0$

$S^3 \rightarrow S^3$ integer topological invariant

Topological solitons: Baryons

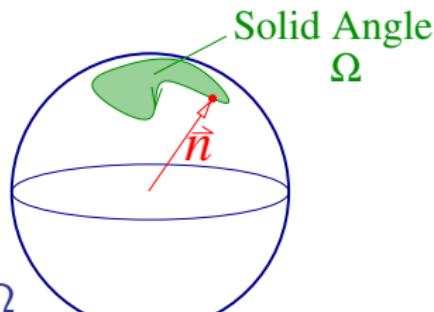
[Battye & Sutcliffe, PRL **79**, 363 (1997)]

“Baby” Skyrmions: $S^2 \rightarrow S^2$

$\vec{n}(\mathbf{r}) = (n_1, n_2, n_3)$ with $|\vec{n}|^2 = 1$
in 2D, with $\vec{n}(\mathbf{r} \rightarrow \infty) = \vec{n}_0$

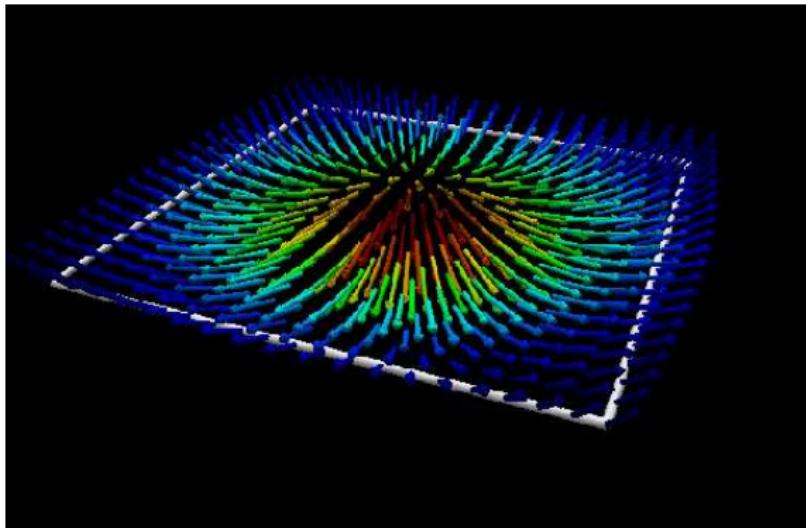
Topological density

$$q(\mathbf{r}) = \frac{1}{8\pi} \epsilon_{ijk} \epsilon_{\mu\nu} n_i \nabla_\mu n_j \nabla_\nu n_k$$



$$\int_A q(\mathbf{r}) d^2\mathbf{r} = \frac{\Omega}{4\pi}$$

Skyrmion: $S^2 \rightarrow S^2$



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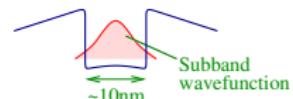
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Quantum Hall Systems

Two-dimensional electron gas

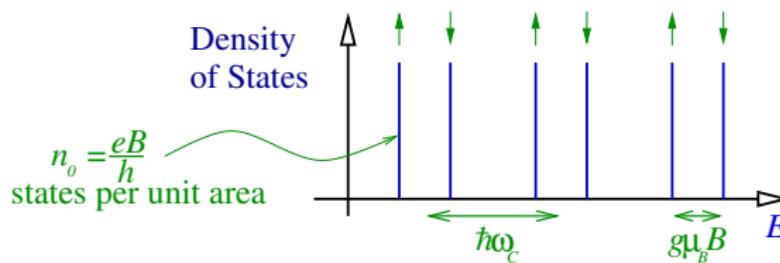


+ strong perpendicular magnetic field B



$$\omega_c = \frac{eB}{m}$$

$$E = \hbar\omega_c(n_l + 1/2)$$



Degeneracy

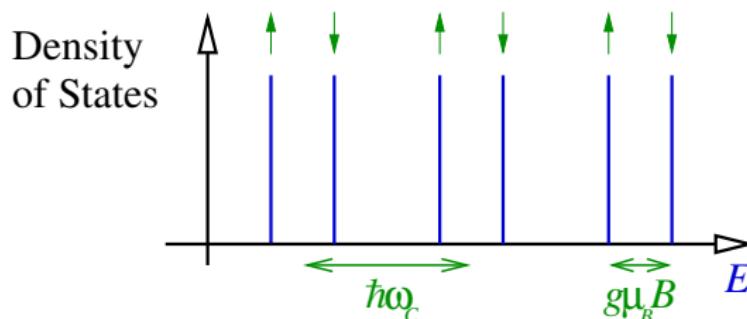
$$N_\phi = \frac{BA}{h/e}$$

Filling factor

$$\nu \equiv \frac{N_e}{N_\phi}$$

QH Ferromagnet at $\nu = 1$

$N_e = N_\phi$: one filled Landau level for one spin.



Ferromagnet \Rightarrow groundstate is fully polarized even for $g\mu_B B \rightarrow 0$.

Low-energy excitations described by local magnetic orientation $\vec{n}(\mathbf{r})$

Topological density \iff electric charge density

Geometric Phases for Itinerant Electrons



Aharonov-Bohm phase

$$\phi_{AB} = 2\pi \frac{e}{h} \int_A B \, d^2r$$

$$B^* = B + \frac{\hbar}{e} q(\mathbf{r})$$

Berry phase

$$\phi_{Berry} = \frac{\Omega}{2} = 2\pi \int_A q(\mathbf{r}) \, d^2r$$

Skyrmiion: $\int (B^* - B) \, d^2r = \frac{\hbar}{e} \times 1 \Rightarrow \text{extra charge } -e$

Effective Action

[Sondhi *et al.*, Phys. Rev. B **47**, 16419 (1993)]

$$\begin{aligned}
 E \simeq & \frac{1}{2} \rho_s \int (\nabla_\mu n_i)^2 d^2\mathbf{r} + g\mu_B B n_0 \frac{1}{2} \int (1 - n_3) d^2\mathbf{r} \\
 & + \frac{e^2}{4\pi\epsilon_0} \int \int \frac{q(\mathbf{r})q(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^2\mathbf{r} d^2\mathbf{r}'
 \end{aligned}$$

Scaling $E \sim \rho_s \mathcal{L}^0 + g\mu_B B n_0 \mathcal{L}^2 + \frac{e^2}{4\pi\epsilon_0 \mathcal{L}}$

$$\mathcal{L} \sim \left(\frac{e^2}{\epsilon g\mu_B B n_0} \right)^{1/3} \Rightarrow \text{Skyrmion size diverges for } g\mu_B B \rightarrow 0$$

Charged excitations carry a large spin

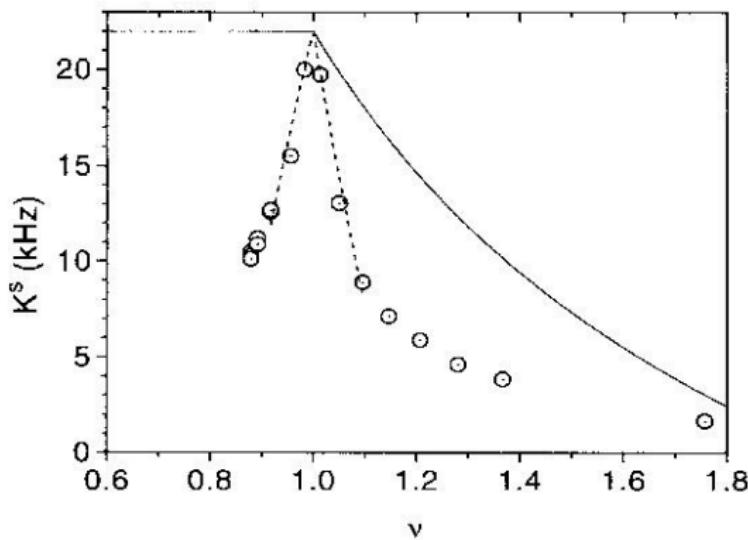
Predicted Experimental Consequences

- ▶ Formation of Skyrmions in the groundstate

$$\nu \equiv \frac{N_e}{N_\phi} = 1 \pm \epsilon$$

- ▶ Rapid loss of spin polarization
- ▶ Enhanced nuclear-spin relaxation rate
- ▶ Skyrmion crystal
- ▶ Dissipation by Skyrmion/antiSkyrmion pairs

Experimental Measurements: Spin Polarization



[Barrett *et al.*, PRL 74, 5112 (1995)]

(anti-)Skyrmions with
average spin 3.6

Fractional Quantum Hall State at $\nu = 5/2$

$$\frac{5}{2} = \underbrace{2}_{\text{filled } n_L = 0 \text{ LL for both spins}} + \underbrace{\frac{1}{2}}_{\text{half-filled } n_L = 1 \text{ LL for one spin}}$$

filled $n_L = 0$ LL for both spins half-filled $n_L = 1$ LL for one spin

The “Moore-Read” state: [Moore & Read, Nucl. Phys. B360, 362 (1991)]

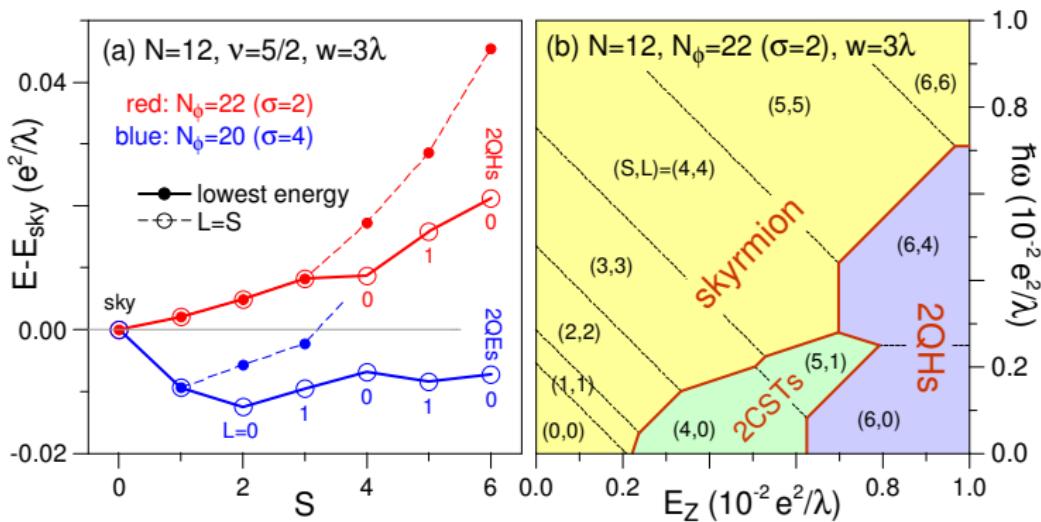
$e/4$ excitations should obey *non-abelian* exchange statistics.

$e/4$ excitations bind into $e/2$ Skyrmions at small Zeeman energies.

[A. Wöjs, G. Möller, S. Simon, and NRC, Phys. Rev. Lett. 104, 086801 (2010)]

Fractional Quantum Hall State at $\nu = 5/2$

[A. W  js, G. M  ller, S. Simon, and NRC, Phys. Rev. Lett. 104, 086801 (2010)]



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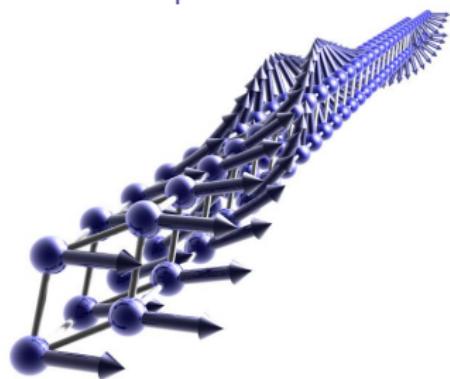
Chiral Magnets

Dzyaloshinskii-Moriya interaction

$$H = \int d^3\mathbf{r} \left[\frac{1}{2} \rho_s (\nabla_i n_j)^2 + \alpha \epsilon_{ijk} n_i \nabla_j n_k \right] + \dots$$

Spin spiral $\vec{n} = (0, \cos qx, \sin qx)$

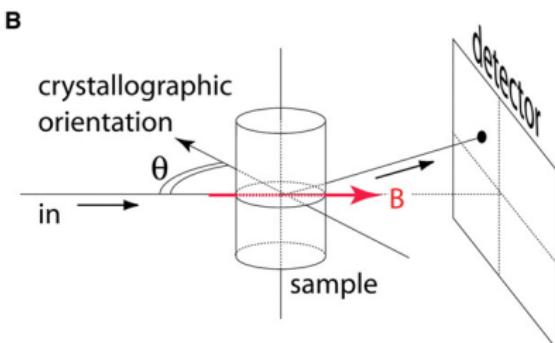
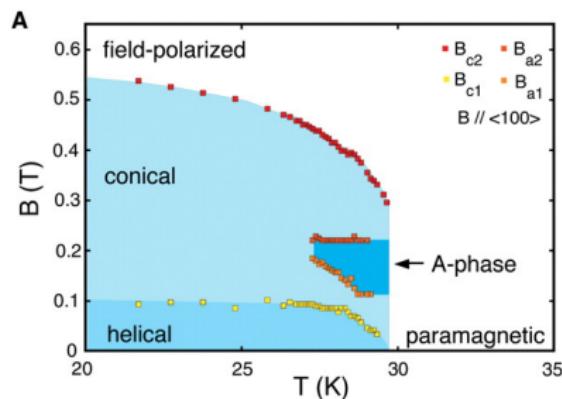
$$q = \frac{\alpha}{\rho_s}$$



MnSi: Phase Diagram

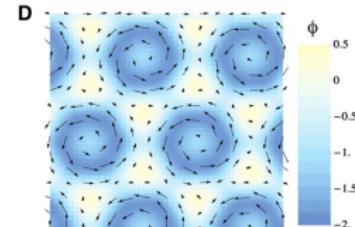
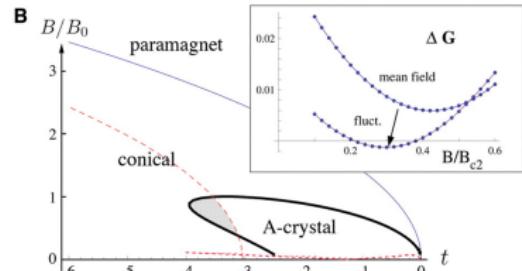
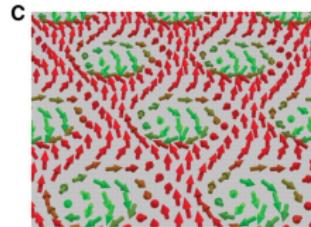
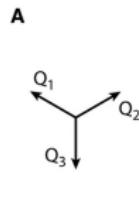
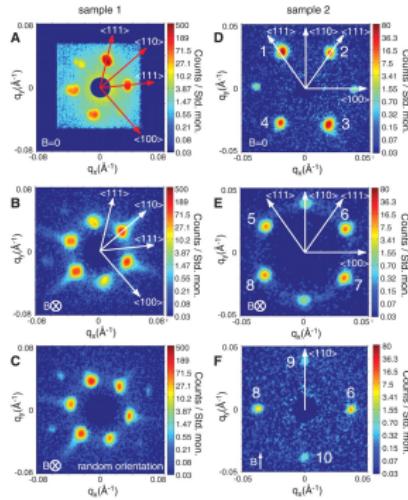
- + Zeeman field $g\mu_B B \frac{1}{2}(1 - n_3)$
- + temperature T

[Mühlbauer *et al.*, Science 323, 915919 (2009)]

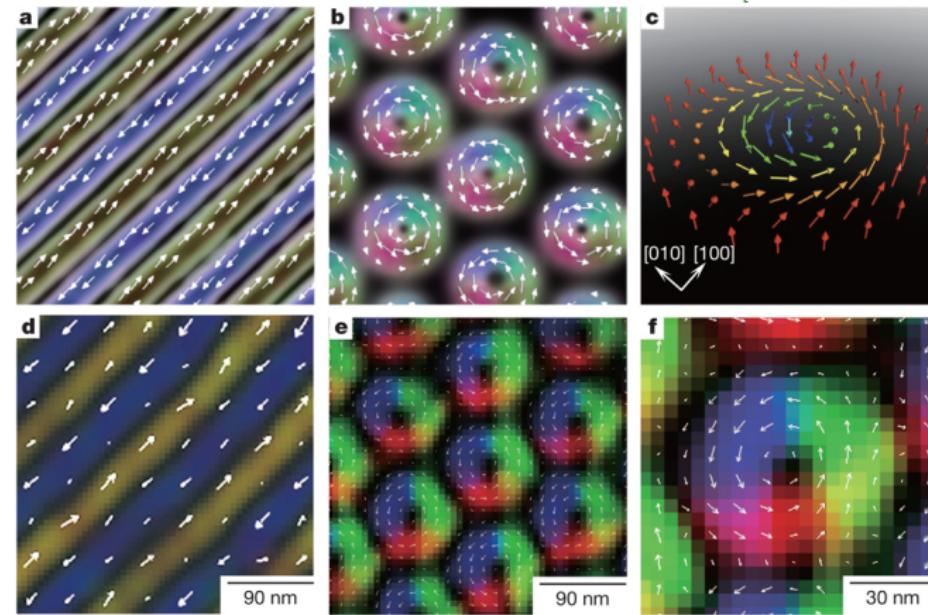


MnSi: Neutron Scattering

[Mühlbauer *et al.*, Science 323, 915919 (2009)]



$\text{Fe}_{1-x}\text{Co}_x\text{Si}$: Real Space Images



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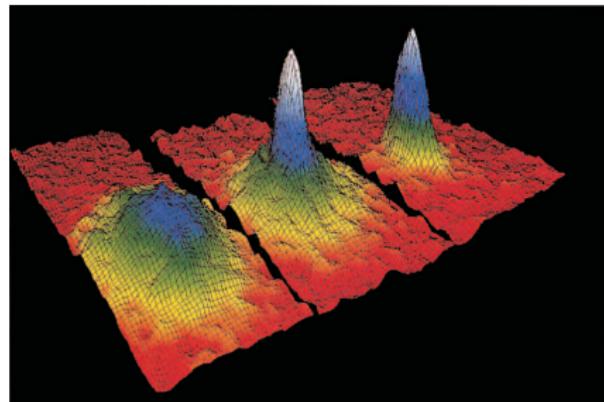
Chiral Magnets

Atomic BECs

Summary

Atomic Bose-Einstein Condensates

[M. H. Anderson *et al.*, Science 269, 198 (1995)]



$$k_B T_c \simeq \frac{\hbar^2}{m} n^{2/3} \sim 100 \text{nK}$$

$$\Psi_N \propto \prod_{i=1}^N \psi(\mathbf{r}_i)$$

$$E = \int d^3\mathbf{r} \left[\frac{\hbar^2}{2m} |\nabla\psi|^2 + V(\mathbf{r})|\psi|^2 + \frac{1}{2} U|\psi|^4 \right]$$

Spinor BECs

Hyperfine states, α $\Psi_N \propto \prod_{i=1}^N \left[\sum_{\alpha} \psi_{\alpha}(\mathbf{r}_i) |\alpha_i\rangle \right]$

$$E = \int d^3\mathbf{r} \left[\sum_{\alpha} \frac{\hbar^2}{2m} |\nabla \psi_{\alpha}|^2 + V_{\alpha}(\mathbf{r}) |\psi_{\alpha}|^2 + \frac{1}{2} \sum_{\alpha, \beta} U_{\alpha\beta} |\psi_{\alpha}|^2 |\psi_{\beta}|^2 \right]$$

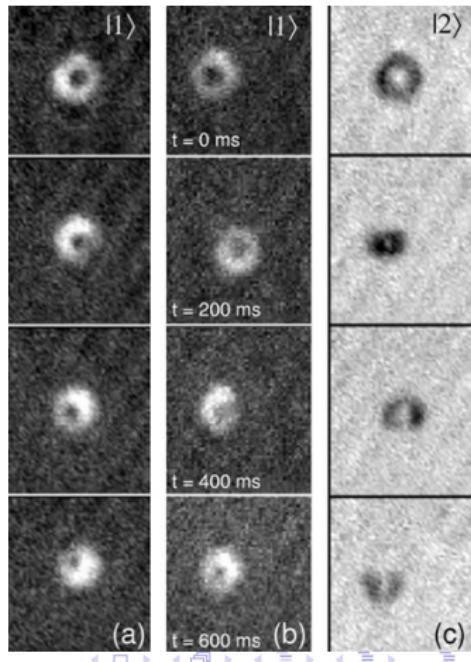
$$\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \vdots \end{pmatrix} \quad \Rightarrow \quad \text{topological solitons ?}$$

Spinor BECs: 2D Skyrmions

[Matthews *et al.*, PRL **83**, 2498 (1998)]

$$\Psi(r) = \sqrt{n} \begin{pmatrix} \cos(\theta/2) \\ e^{i\varphi} \sin(\theta/2) \end{pmatrix}$$

$$\langle \vec{\sigma} \rangle = (\sin \theta \cos \varphi, \sin \theta \sin \varphi, \cos \theta)$$

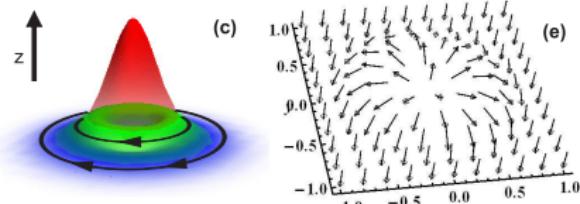
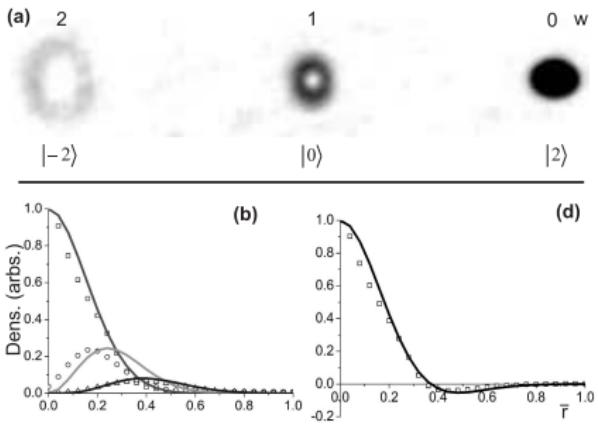


Spinor BECs: 2D Skyrmions

[Leslie *et al.*, PRL 103, 250401 (2009)]

$$\sqrt{n} \begin{pmatrix} \cos^2(\theta/2) & & \\ 0 & & \\ \sqrt{2}e^{i\varphi} \sin(\theta/2) \cos(\theta/2) & & \\ 0 & & \\ e^{2i\varphi} \sin^2(\theta/2) & & \end{pmatrix}$$

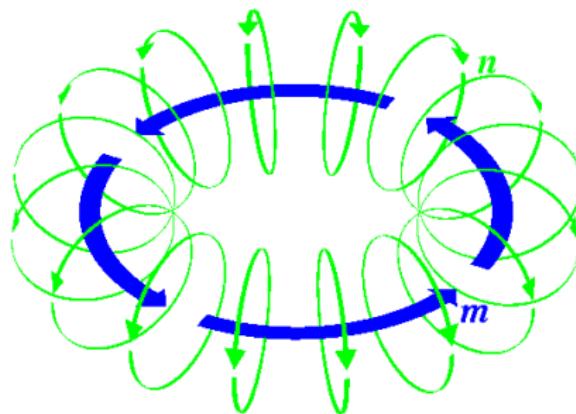
$$\vec{\ell} = (\sin \theta \cos \varphi, \sin \theta \sin \varphi, \cos \theta)$$



Spinor BECs: 3D Skyrmions

Two-component BEC

$$\begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} = \sqrt{n} \begin{pmatrix} n_1 + i n_2 \\ n_3 + i n_4 \end{pmatrix}$$
$$n_1^2 + n_2^2 + n_3^2 + n_4^2 = 1$$



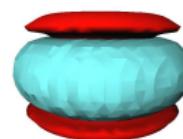
$$Q = mn$$

[cf. "Cosmic vortons"]

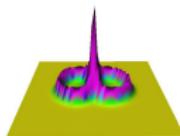
Stable 3D Skyrmions

[R. Battye, NRC & P. Sutcliffe, PRL (2002)]

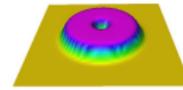
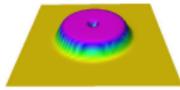
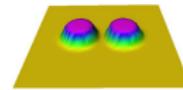
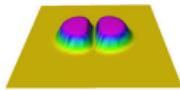
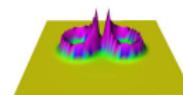
$$U_{12}^2 > U_{11} U_{22}$$



$Q = 1$



$Q = 2$



Summary

Skyrmions appear in a variety of condensed matter settings.

- ▶ Quantum Hall Ferromagnets:
topological density \Leftrightarrow electric charge density
- ▶ Chiral Magnets: Skyrmion crystal phase
- ▶ Spinor BECs