

Lecture 3: Matrix-Product states.

- idea:
- use representation of eigenstates that focuses on the entanglement.
 - typical physical states have mostly local entanglement, so using the full basis of the Hilbert-space is not necessary!
 - want to use this to circumvent exponential growth of complexity in N .

Let us first think about spin-systems, again:

$$\begin{array}{ccccccc}
 s_1 & & & & & i & s_N \\
 \uparrow & \uparrow & \uparrow & \uparrow & \dots & \uparrow & \uparrow \\
 \uparrow & \uparrow & \uparrow & \uparrow & \dots & \uparrow & \uparrow \\
 & & & & \dots & & \\
 & & & & \text{local Hilbert-space } V_n = \{|\uparrow\rangle, |\downarrow\rangle\} & & \\
 \text{total Hilbert-space } \mathcal{H} = \bigotimes_{n=1}^N V_n & & & & & &
 \end{array}$$

$$\text{Hamiltonian} = \sum_{\langle i,j \rangle} \hat{h}_{ij} \leftarrow \text{local terms.}$$

general form of the ground state

$$|\psi\rangle = \sum_{\substack{s_1 \dots s_N \\ s_k \in \{\uparrow, \downarrow\}}} \psi_{s_1 s_2 s_3 \dots s_N} |s_1\rangle \otimes |s_2\rangle \otimes \dots \otimes |s_N\rangle$$

$\uparrow$
 $2^N \text{ coefficients} = \text{total Hilbert-space}$

Here, read $\psi_{s_1 \dots s_N}$ as a matrix / tensor with N indices $s_1 \dots s_N$.

Let's learn to write these efficiently

(2)

Diagrammatic notation

- each tensor is represented by a box / shape $\bigcirc$
- external legs represent indices

$$\bigcirc = \text{C-number } c$$

$$\bigcirc - i = \text{vector } v_i$$

$$i - \bigcirc - j = \text{matrix } M_{ij}$$

$$i - \bigcirc - j - k = \text{rank 3 Tensor } T_{ijk}$$

⋮

- Summation-convention: lines connecting different tensors correspond to indices to be summed over

$$i - \bigcirc M - j - \bigcirc V = i - \bigcirc u$$

$$\sum_j M_{ij} v_j = u_i$$

or

$$\bigcirc u - i - \bigcirc T - j - \bigcirc V = \bigcirc C$$

$\bigcirc$
 k
 $\bigcirc w$

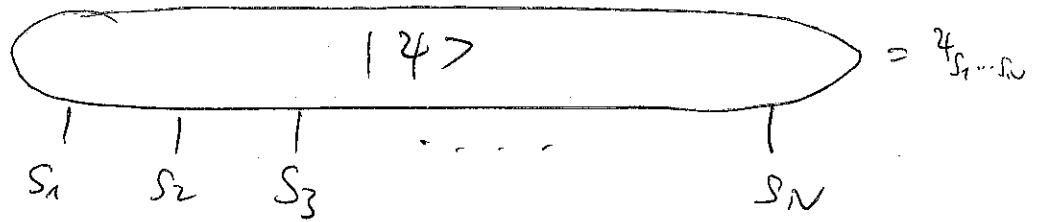
or

$$\begin{array}{c} \bigcirc \\ \diagup \quad \diagdown \\ \bigcirc \quad \bigcirc \\ | \quad | \\ r \quad s \end{array} = i - \bigcirc M - s$$

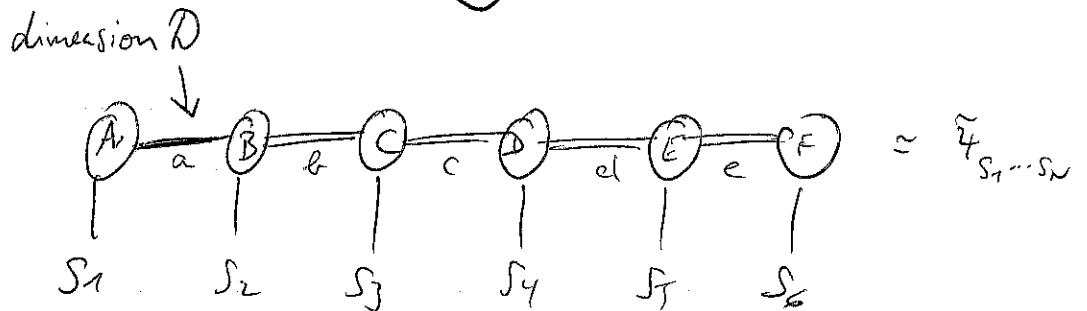
(Show big example on slide)

Representing the wave-function : spin systems

2^N coefficients $\Rightarrow$
 $\sim \exp(N)$



$\Downarrow$ Ansatz

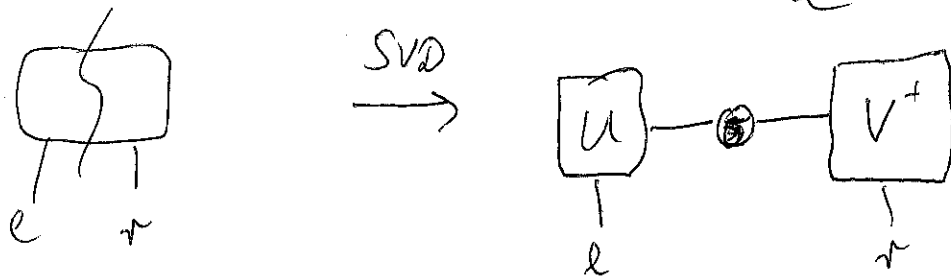
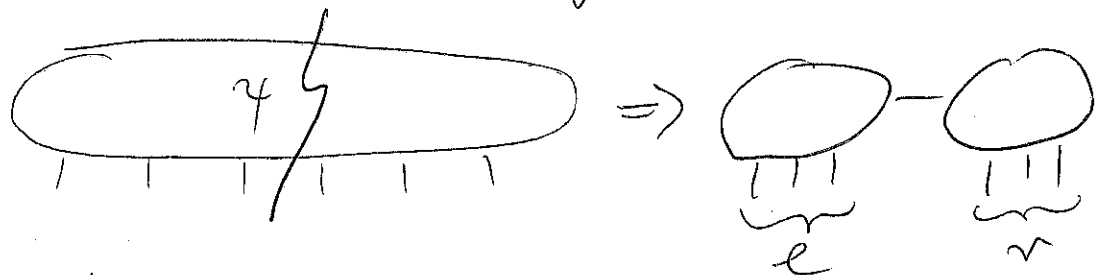


$$= \sum_{a \dots e} A_{s_1}^a B_{s_2}^{ab} C_{s_3}^{bc} \dots F_{s_6}^e$$

parameters $\approx 2 \cdot N \cdot D^2 \approx \text{poly}(N, D)$.

Why would such an Ansatz work?

let's think about splitting tensors



$$\begin{aligned} \psi_{er} &= \sum_{l, k} \psi_{er} |l\rangle |r\rangle = \sum_{e, r} \sum_k U_{el} V_{rk}^\dagger |l\rangle |r\rangle \\ &= \sum_k \sigma_{rk} |u_k\rangle |v_k\rangle \end{aligned}$$

(4)

Examples: pair of spins

• Spin-triplet

$$|4\rangle = |\uparrow\rangle_1 \otimes |\uparrow\rangle_2 = S_z \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}^{S_1}$$

$$= \sum_{\sigma} S_{\sigma} |\uparrow\rangle_1 |\sigma\rangle_1 \quad \text{with } S_+ = 1, S_- = 0$$

unentangled state, only one singular value!

• Spin-singlet

$$|4\rangle = \frac{1}{\sqrt{2}} (|\uparrow\rangle_1 \otimes |\downarrow\rangle_2 - |\downarrow\rangle_1 \otimes |\uparrow\rangle_2)$$

$$= \begin{pmatrix} 0 & 1/\sqrt{2} \\ -1/\sqrt{2} & 0 \end{pmatrix} = \underbrace{\begin{pmatrix} 1 & \\ & 1 \end{pmatrix}}_U \underbrace{\begin{pmatrix} 1/\sqrt{2} & \\ & 1/\sqrt{2} \end{pmatrix}}_{\sigma} \underbrace{\begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}}_{V^\dagger}$$

$$= \frac{1}{\sqrt{2}} |u_1\rangle |v_1\rangle + \frac{1}{\sqrt{2}} |u_2\rangle |v_2\rangle$$

maximally entangled state $\rightarrow$ ~~# singular values = # det.~~
all singular values have equal magnitude

Realistic states:

entanglement grows with the length of the circumference, not volume.

• show example from EO:

reduced density matrix.

Examples of Tensor Networks $\Rightarrow$ continued on slides